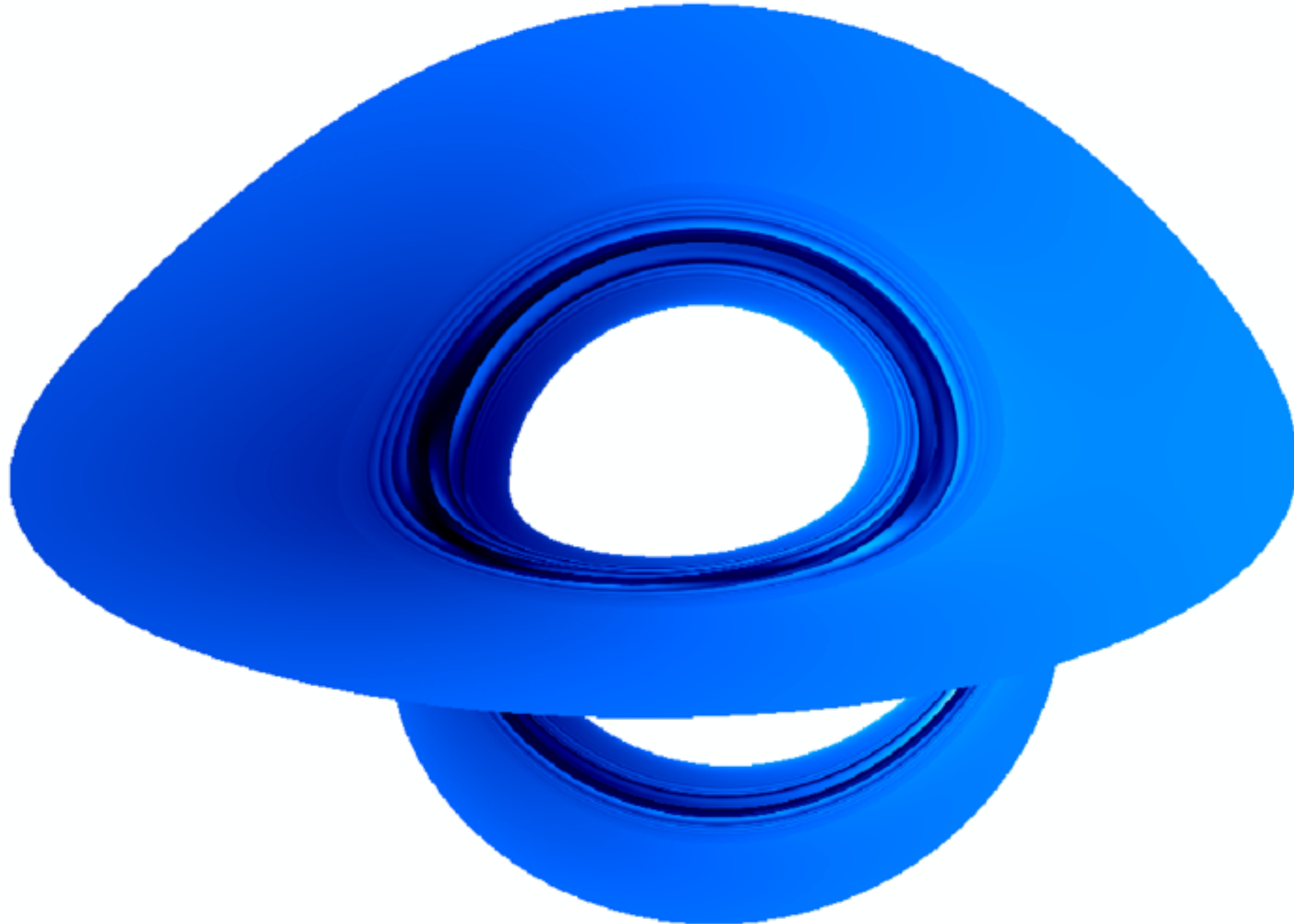


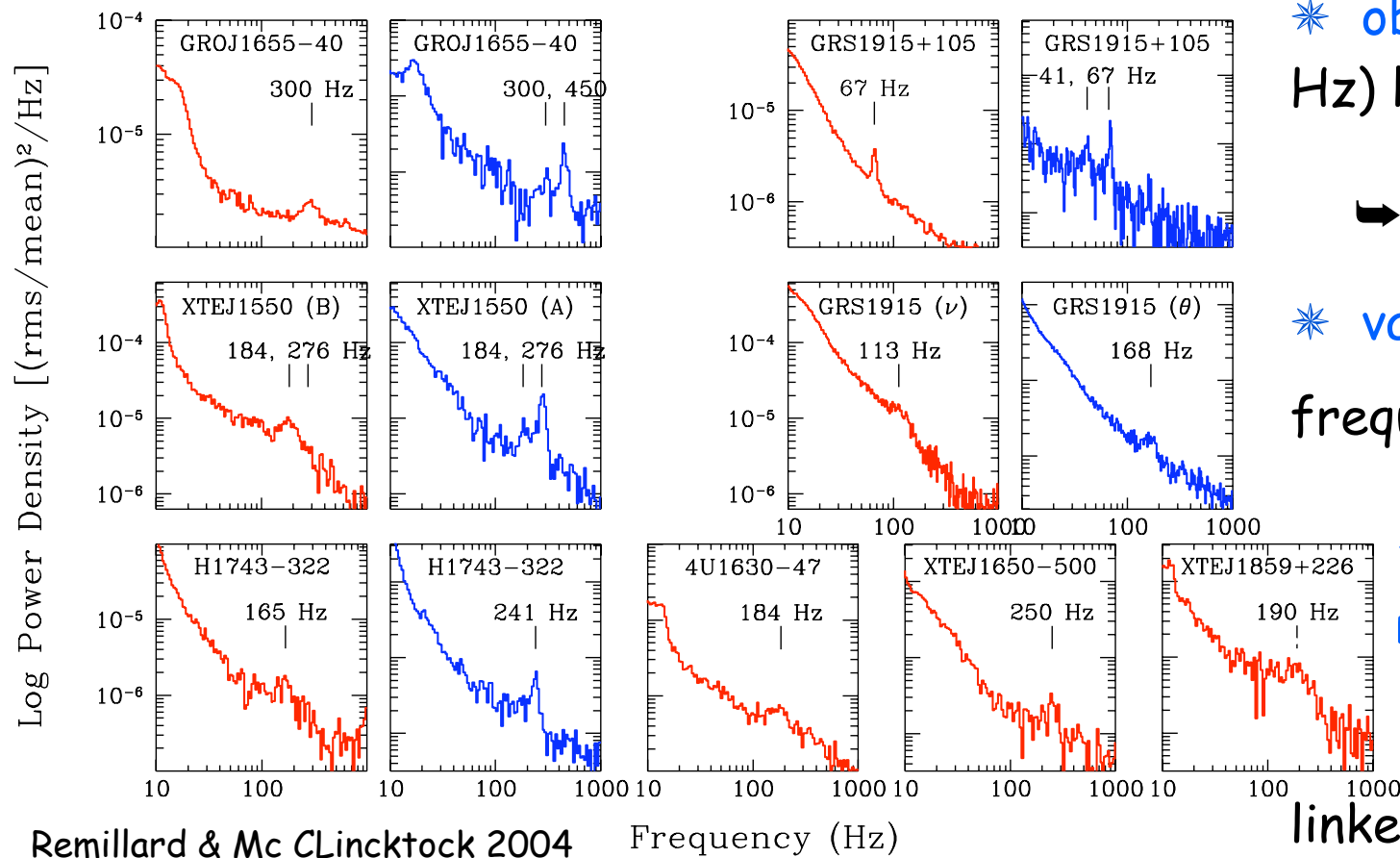
Following the Rossby Wave Instability into the Kerr Metric



P. Varniere, F.Casse, F.Vincent

why the Rossby Wave Instability in full GR?

The **Rossby Wave Instability (RWI)** is one of the few models proposed to explain high frequency Quasi-Periodic Oscillations in black hole systems.



* observed frequency range (up to a few hundreds Hz) link this phenomena with the last stable orbit
 → GR is needed to fully compute the phenomena

* variations in the observed frequency (change in frequency of about +/- 15%)

* HFQPOs appear alone or in "pairs" (with related frequencies)

⇒ need one mechanism that can select several linked frequencies depending on the disk conditions

* HFQPOs can occur in the absence of LFQPOs

* when they co-exist we have "unusual" LFQPO (type A and B)

* HFQPOs rms amplitudes are much lower than LFQPO and seem anti-correlated with the LFQPO rms

⇒ the HFQPO model need to be **coherent with a LFQPO model** as they have to co-exist in the disk while being independent

⇒ the HFQPO model also need to be able to develop in a **wide range of disk conditions** as LFQPOs

can reach high rms value hence changing the condition in which the HFQPO model exists

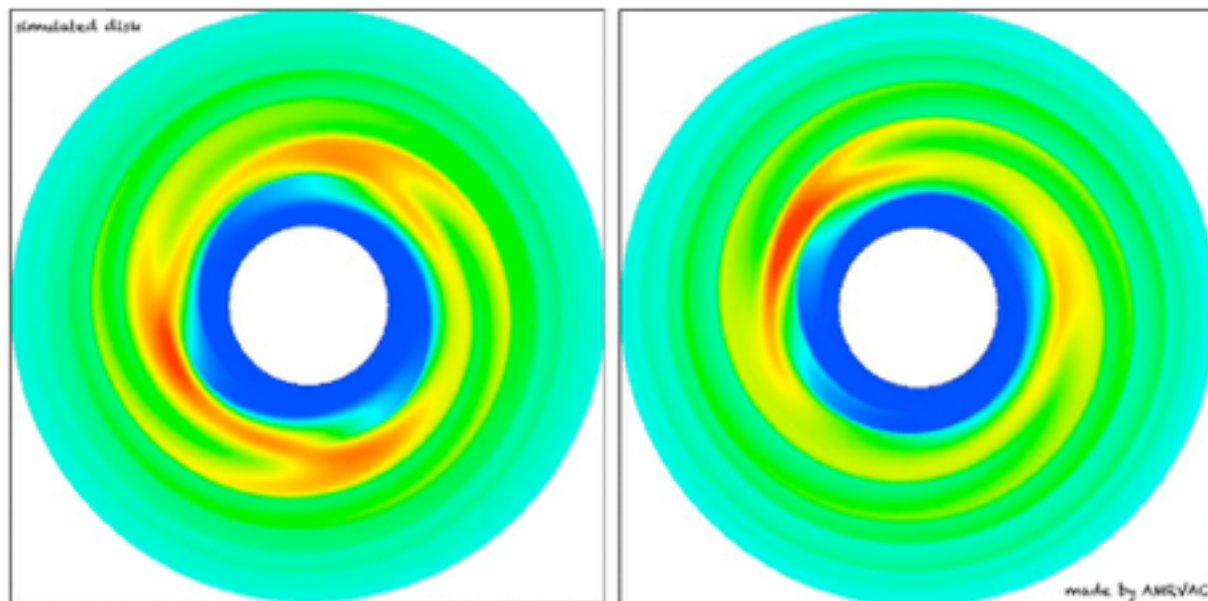
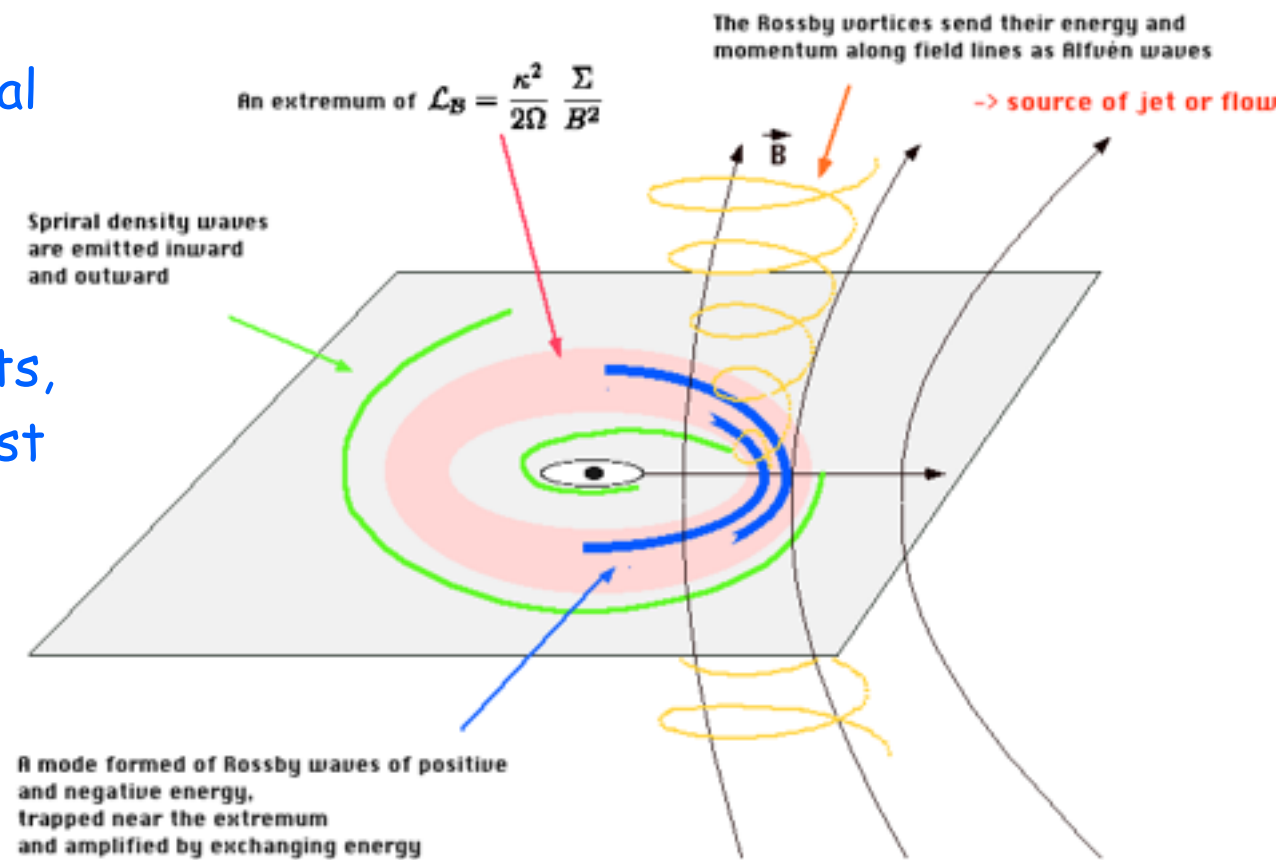
classical RWI

The Rossby Wave Instability (RWI) is a hydrodynamical instability proposed to explain high frequency Quasi-Periodic Oscillations in black hole systems.

it requires having an extremum of vortensity which exists, for example, in disks with their inner edge close to the last stable orbit.

$$\mathcal{L} = \frac{(\nabla \times \mathbf{v})|_z}{\Sigma} \text{ or } \mathcal{L}_B = (\nabla \times \mathbf{v})|_z \cdot \frac{\Sigma}{B^2}$$

→ the existence of the RWI is linked to the position of the vortensity extremum



* One interesting characteristic of this instability is that, depending on physical condition in the disk, the $m=1$ mode is not dominant but it is rather a mix of the $m=2, 3$ and higher modes that dominates.

→ here we have one mechanism that can select several linked frequencies depending on the disk conditions

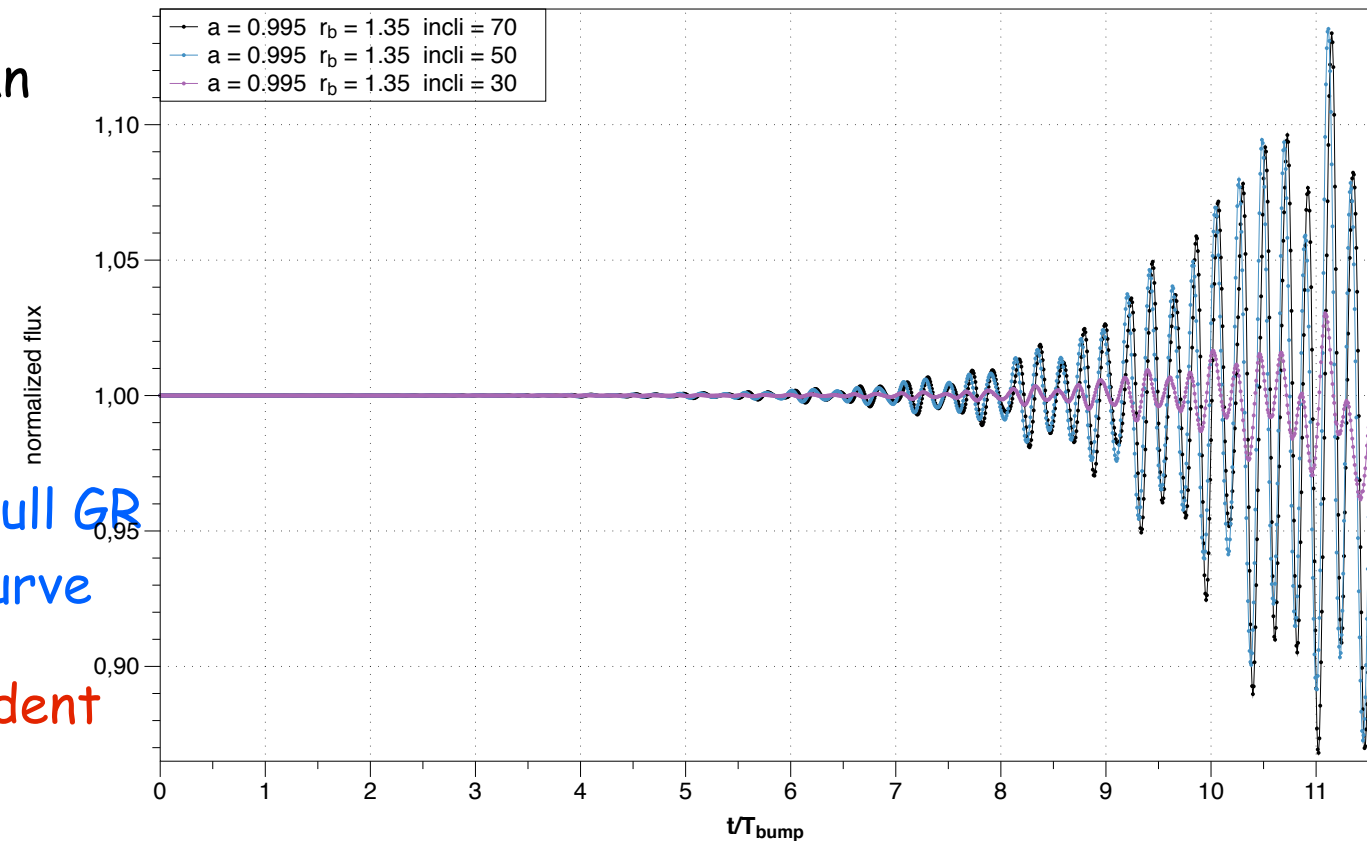
RWI from PN to GR

at first we studied the RWI in the Pseudo-Newtonian approach with $\mathcal{L} = \frac{\kappa^2}{2\Omega\Sigma}$

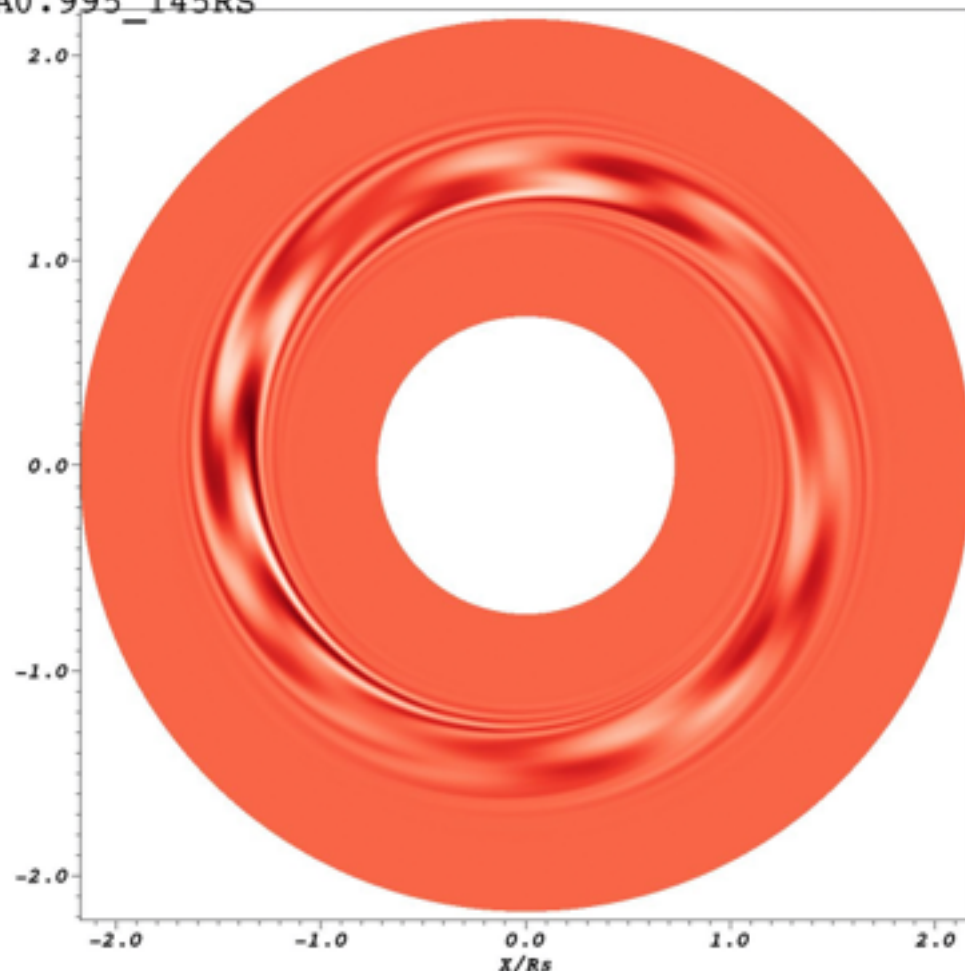
and were able to show that the RWI is triggered by the maximum of the epicyclic frequency

* using 3D simulation of this instability coupled with full GR ray-tracing we computed the associated image/light curve

⇒ a few % modulation is visible and is energy dependent



DB: A0.995_145RS

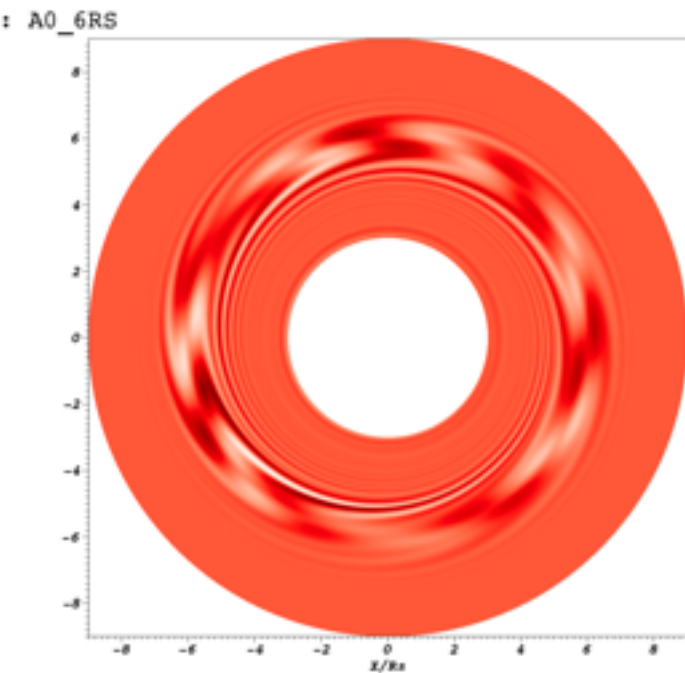


* using the newly developed GR-AMRVAC code we were able to follow the RWI from a purely Newtonian setup (at 150 rg) to a case with $a=0.995$ and the inner edge of the disk at its last stable orbits.

→ hence proving that the RWI exist and is unstable in a full GR calculation.

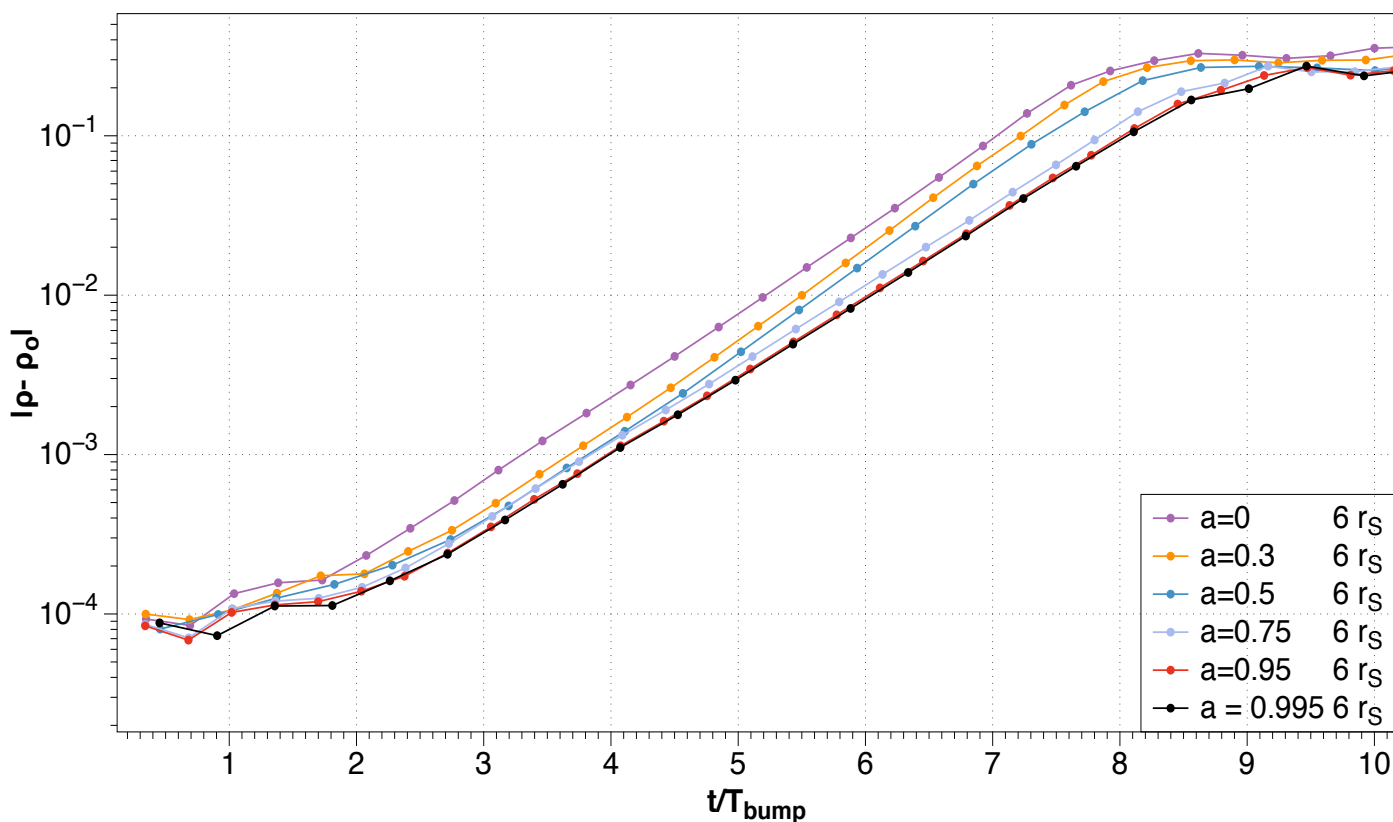
Increasing the spin

spin from 0 to 0.995



instability from 12 r_g to 2.90 r_g

increasing the spin and getting closer to the black hole **the instability looks qualitatively the same**



looking at the same location/frequency, for different spin **we see a small difference in the growth time of about one orbits and a relatively similar saturation level**

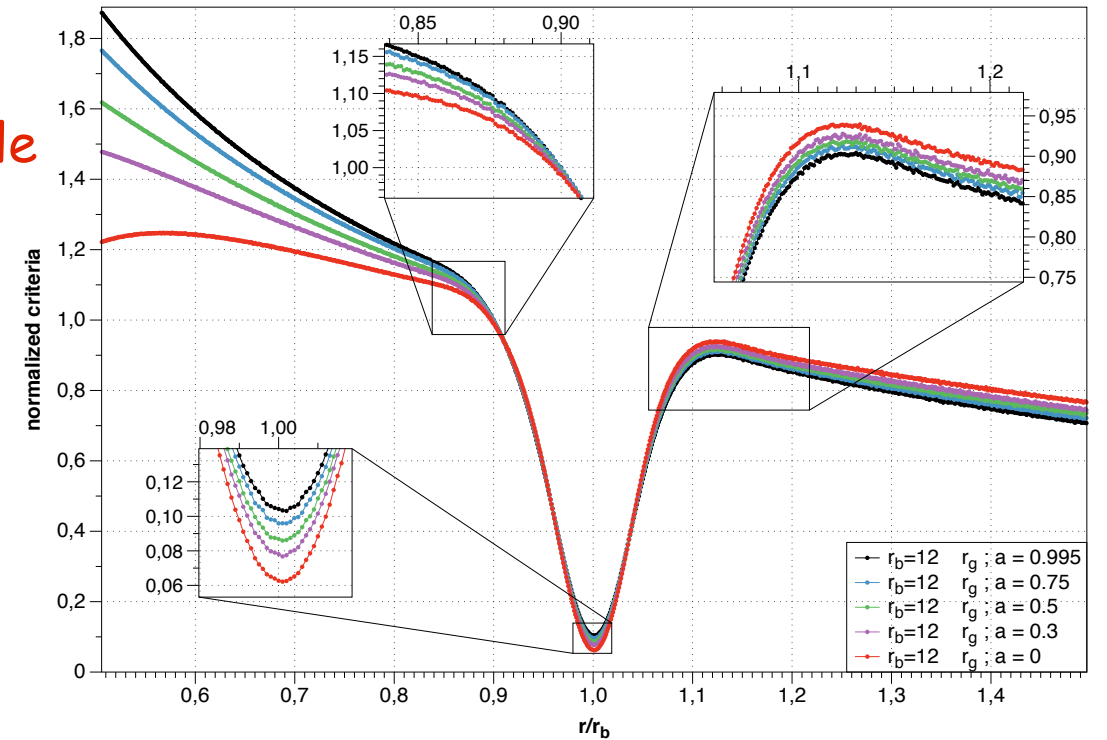
=> no direct impact of spin at a given frequency on the instability behavior

=> but the perception the observer has of the instability has to take into account GR

plunging into the Kerr Black Hole

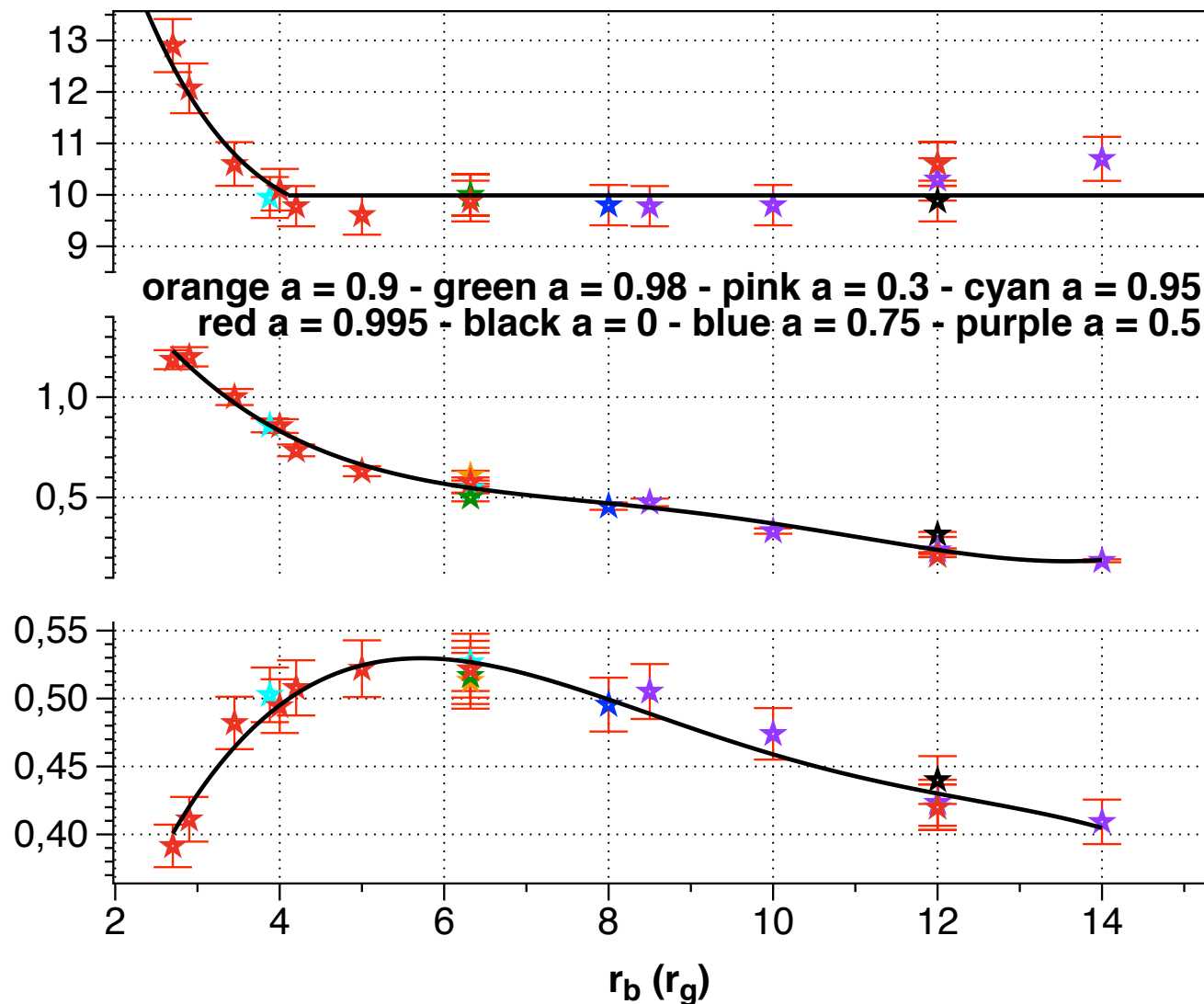
* As the spin increases, we can reach closer to the black hole giving access to higher frequencies

➔ and the RWI will develop in a stronger gravity well



while no impact was detectable at a given frequency we see that as the spin increases:

- the growth rate seems to reach a maximum then decreases
- the RWI reach a higher saturation but on a slightly longer timescale



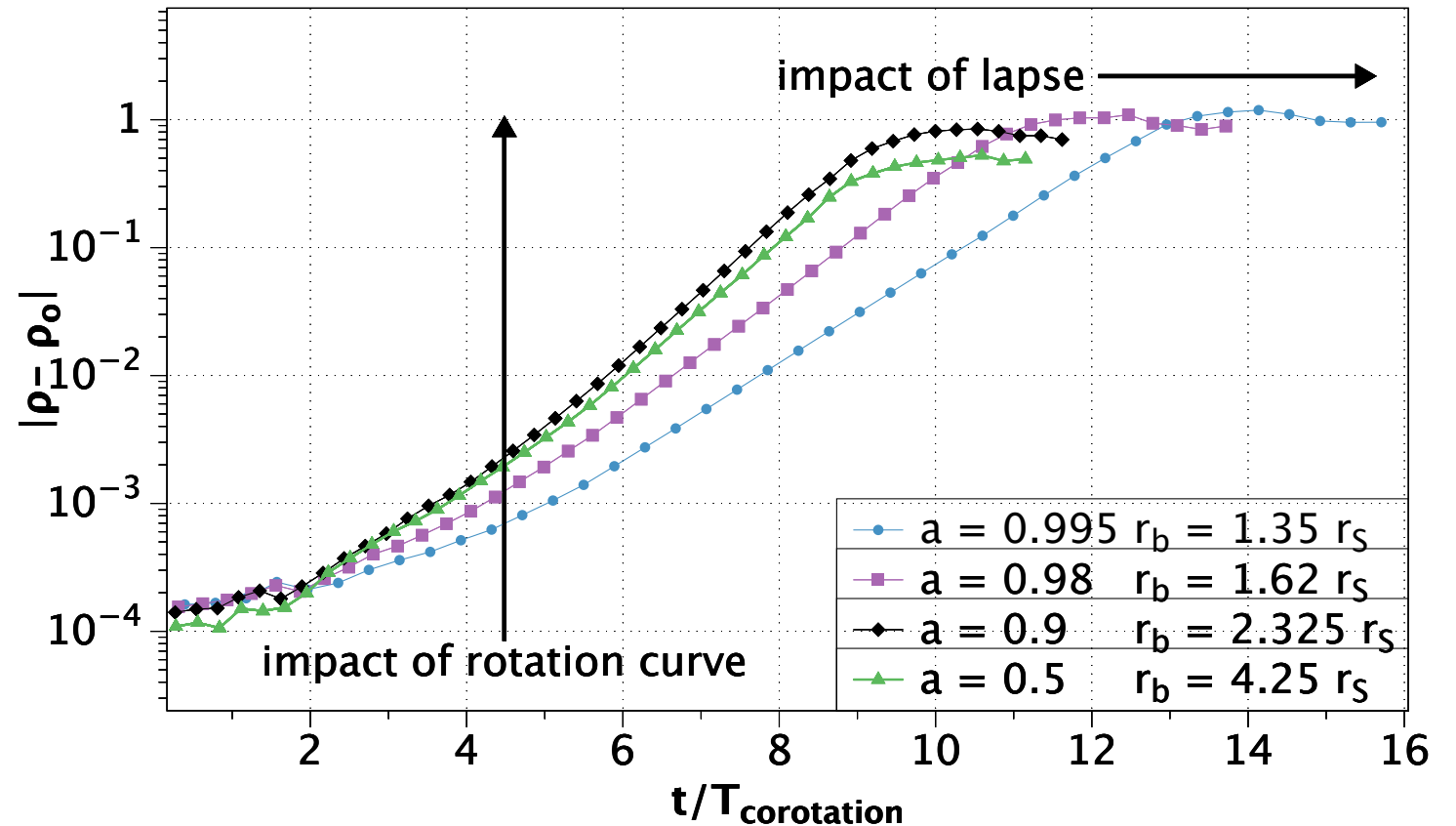
* the next step is to focus on observable as we did in the PN case by computing both energy and power density spectrum

⇒ We are also working on better understanding the constraints from observations.

link with observables

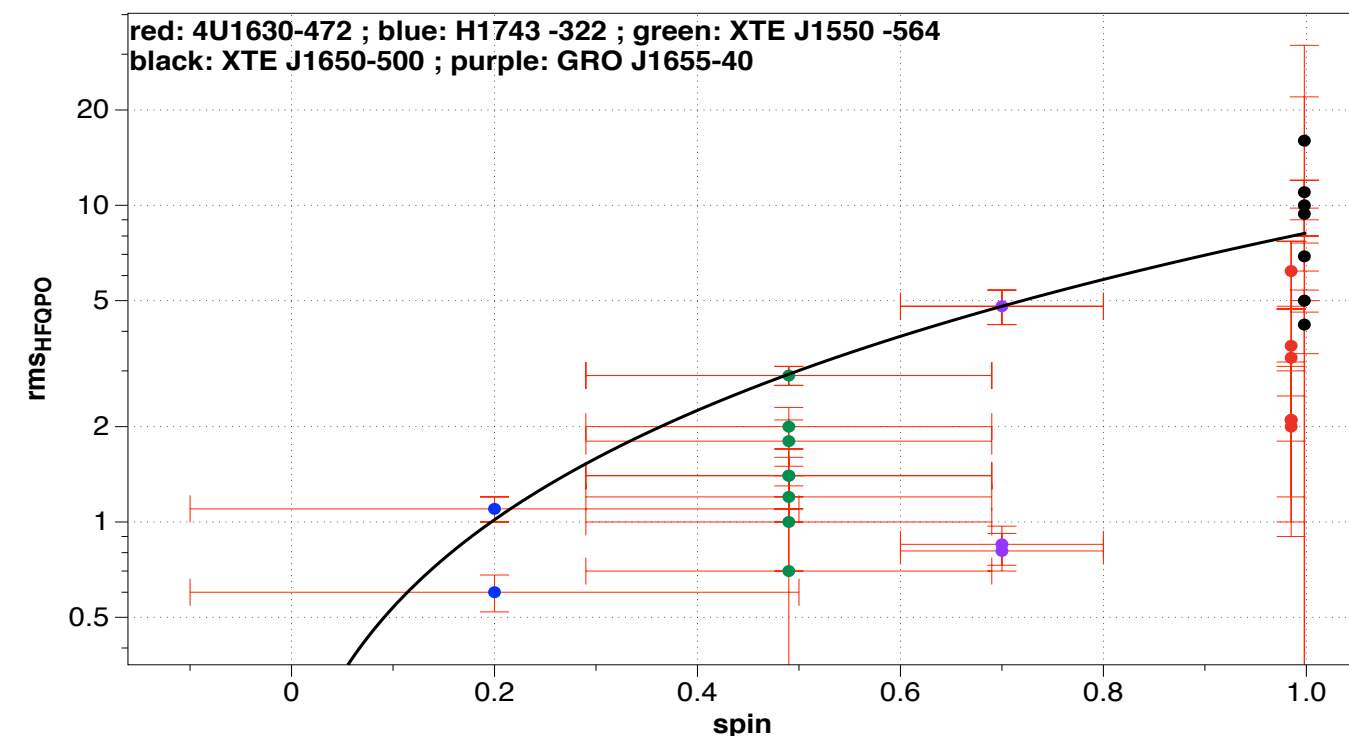
* in the HFQPO model the RWI will be triggered when the inner edge of the disk gets close to the last stable orbit of the black hole

→ we can compare the behaviour of the RWI when for different LSO/spin couple



the main difference is that higher spin implies a higher saturation level for the instability

while this is not a direct observable, we can look at the maximum amplitude observed for each black hole with a known spin



* with only a few HFQPO sources with known spin, and taking into account the spin controversy, it is a weak observable

⇒ getting more spin measurement for the rest of the HFQPO sources would help

→ we are also looking for other observables

Conclusion

- ⌘ The RWI is an instability that can occur in the strong gravity of a Kerr black hole as it was theorized.
- ⌘ The spin of the black hole does have an impact on the instability and through it to the detectability of HFQPO
 - but it is indirect as it is related to how close to the black hole the accretion disk, and hence where the instability, can developed.
 - nevertheless, it might lead to observational tests and we are looking for more spin/HFQPO couples.
- ⌘ We are also working on creating synthetic observations related to our simulations, and also an XSPEC model to directly fit observation.
 - ⇒ directly comparing the impact on the energy spectra and PDS